

Supplementary Materials for “Modeling Tangential Vector Fields on a Sphere”

S.1 Additional Technical Details

S.1.1 Auxiliary Notations and Definitions

In this subsection, we describe some auxiliary notations and definitions helpful for understanding the construction and characteristics of tangential vector fields on the unit sphere.

The matrix \mathbf{Q}_s is defined as

$$\mathbf{Q}_s =: \hat{\mathbf{r}} \times \mathbf{P}_s = \begin{pmatrix} 0 & -s_3 & s_2 \\ s_3 & 0 & -s_1 \\ -s_2 & s_1 & 0 \end{pmatrix}. \quad (\text{S.1})$$

The surface gradient and the surface curl operators can be represented in spherical coordinates except at the two poles, i.e.,

$$\nabla_s^* f = \frac{\partial f}{\partial \theta} \hat{\boldsymbol{\theta}} + \frac{1}{\sin \theta} \frac{\partial f}{\partial \phi} \hat{\boldsymbol{\phi}},$$

and

$$L_s^* f = \frac{\partial f}{\partial \theta} \hat{\boldsymbol{\phi}} - \frac{1}{\sin \theta} \frac{\partial f}{\partial \phi} \hat{\boldsymbol{\theta}}.$$

Let $X(\mathbf{s})$ denote a scalar random field on a spherical shell S_ϵ , where $(\Omega, \mathcal{F}, \mathbb{P})$ is the underlying probability space. The partial derivative of $X(\mathbf{s})$ in quadratic mean along the i -th coordinate

direction is defined as a stochastic process $D_{\text{qm}}^{(i)}X(\mathbf{s})$ such that

$$\mathbb{E} \left(\frac{X(\mathbf{s} + h\mathbf{e}_i) - X(\mathbf{s})}{h} - D_{\text{qm}}^{(i)}X(\mathbf{s}) \right)^2 \rightarrow 0 \quad \text{as } h \rightarrow 0,$$

where \mathbf{e}_i is the unit vector in \mathbb{R}^3 with the i -th element one and zeros elsewhere. The sample partial derivatives of $X(\mathbf{s})$, $D^{(i)}X(\mathbf{s})$, $i = 1, 2, 3$, are defined as the usual partial derivatives of $X(\mathbf{s}, \omega)$ for any fixed $\omega \in \Omega$. A random field \tilde{X} is said to be a \mathbb{P} -a.e. sample continuously differentiable version of X if the following hold:

- (i) For any $\mathbf{s} \in S_\epsilon$, $\mathbb{P} \left(\tilde{X}(\mathbf{s}) = X(\mathbf{s}) \right) = 1$.
- (ii) There exists $\Omega_0 \subset \Omega$ with $\mathbb{P}(\Omega_0) = 1$ such that $\tilde{X}(\mathbf{s}, \omega)$ is continuously differentiable on S_ϵ for any $\omega \in \Omega_0$.

A random tangential vector field on \mathbb{S}^2 is said to be curl-free (or divergence-free) if its sample paths are curl-free (or divergence-free) \mathbb{P} -a.e.

S.1.2 Proof of Theorem 1

We first verify (1). Since $D^{(i)}\tilde{Z}(\mathbf{s}) = D_{\text{qm}}^{(i)}Z(\mathbf{s})$ \mathbb{P} -a.e., $\mathbf{Y}_{\text{curl},Z}(\mathbf{s}) = \mathbf{P}_s \nabla_{\text{s,qm}} Z(\mathbf{s})$ \mathbb{P} -a.e., where $\nabla_{\text{s,qm}}$ is the usual gradient on \mathbb{R}^3 defined in the sense of quadratic mean. Thus,

$$\begin{aligned} \mathbf{C}_{\text{curl},Z}(\mathbf{s}, \mathbf{t}) &= \mathbf{P}_s \text{Cov} (\nabla_{\text{s,qm}} Z(\mathbf{s}), \nabla_{\text{t,qm}} Z(\mathbf{t})) \mathbf{P}_t^T \\ &= \mathbf{P}_s \left[\text{Cov} (D_{\text{qm}}^{(i)}Z(\mathbf{s}), D_{\text{qm}}^{(j)}Z(\mathbf{t})) \right]_{1 \leq i, j \leq 3} \mathbf{P}_t^T \\ &= \mathbf{P}_s \left[\frac{\partial^2 C}{\partial s_i \partial t_j}(\mathbf{s} - \mathbf{t}) \right]_{1 \leq i, j \leq 3} \mathbf{P}_t^T \quad (\text{by condition A2}) \\ &= -\mathbf{P}_s \left[\frac{\partial^2 C}{\partial h_i \partial h_j}(\mathbf{h}) \right]_{1 \leq i, j \leq 3} \Big|_{\mathbf{h}=\mathbf{s}-\mathbf{t}} \mathbf{P}_t^T \\ &= -\mathbf{P}_s \nabla_{\mathbf{h}} \nabla_{\mathbf{h}}^T C(\mathbf{h}) \Big|_{\mathbf{h}=\mathbf{s}-\mathbf{t}} \mathbf{P}_t^T. \end{aligned}$$

(2) can be derived similarly.

S.1.3 Proof That $\nu > 1$ Is Necessary and Sufficient

First, $\nu > 1$ is necessary because the underlying potential field $Z(\mathbf{s})$ is required to be differentiable in quadratic mean. Second, we shall show the sufficiency. Since $\nu > 1$ and Z is Gaussian, the sample paths of Z are continuously differentiable \mathbb{P} -a.e. (Handcock and Wallis, 1994). Besides, Z has partial derivatives in quadratic mean. These two facts together imply that $(Z(\mathbf{s} + h\mathbf{e}_i) - Z(\mathbf{s})) / h$ converges to $D_{\text{qm}}^{(i)}Z(\mathbf{s})$ in L_2 (and in probability), and converges to $D^{(i)}Z(\mathbf{s})$ almost everywhere (and in probability). By the uniqueness of convergence in probability, $D_{\text{qm}}^{(i)}Z(\mathbf{s}) = D^{(i)}Z(\mathbf{s})$ \mathbb{P} -a.e. Thus, we can choose Z as \tilde{Z} so that condition **A1** is satisfied. Jun and Stein (2007) pointed out that the Matérn model (3) is twice continuously differentiable and $M'(0; \nu, a) = 0$. This implies that $C_1(\|\mathbf{h}\|)$ is also twice continuously differentiable when $\|\mathbf{h}\| > 0$. It can be shown by definition that

$$\frac{\partial C_1(\|\mathbf{h}\|)}{\partial h_i} = \begin{cases} C_1'(\|\mathbf{h}\|) \frac{h_i}{\|\mathbf{h}\|} & \text{if } \mathbf{h} \neq \mathbf{0} \\ 0 & \text{if } \mathbf{h} = \mathbf{0}. \end{cases}$$

Based on the first derivatives, we can also show by definition that

$$\frac{\partial^2 C_1(\|\mathbf{h}\|)}{\partial h_i \partial h_j} \Big|_{\mathbf{h}=\mathbf{0}} = \begin{cases} 0 & \text{if } i \neq j \\ C_1'''(0) & \text{if } i = j. \end{cases}$$

This is equivalent to

$$\nabla_{\mathbf{h}} \nabla_{\mathbf{h}}^T C_1(\|\mathbf{h}\|) \Big|_{\mathbf{h}=\mathbf{0}} = F(0) \mathbf{I}_3.$$

Besides, we know that

$$\lim_{\mathbf{h} \rightarrow \mathbf{0}} \nabla_{\mathbf{h}} \nabla_{\mathbf{h}}^T C_1(\|\mathbf{h}\|) \Big|_{\mathbf{h} \neq \mathbf{0}} = \lim_{\mathbf{h} \rightarrow \mathbf{0}} F(\|\mathbf{h}\|) \mathbf{I}_3 + G(\|\mathbf{h}\|) \mathbf{h} \mathbf{h}^T = F(0) \mathbf{I}_3.$$

Thus, $C_1(\|\mathbf{h}\|)$ is twice continuously differentiable at $\mathbf{h} = \mathbf{0}$, i.e., condition **A2** is satisfied.

S.1.4 Tangential Vector Fields Derived from a Full Bivariate Matérn Model

The underlying potential field $\mathbf{Z}(\mathbf{s})$ is assumed to follow a full bivariate Matérn model (Gneiting et al., 2010), which specifies

$$C_{ii}(\|\mathbf{h}\|) = \sigma_i^2 M(\|\mathbf{h}\|; \nu_i, a_i) \quad \text{for } i = 1, 2,$$

and

$$C_{12}(\|\mathbf{h}\|) = C_{21}(\|\mathbf{h}\|) = \rho_{12} \sigma_1 \sigma_2 M(\|\mathbf{h}\|; \nu_{12}, a_{12}).$$

It is valid if and only if

$$\rho_{12}^2 \leq \frac{\Gamma(\nu_1 + \frac{3}{2})}{\Gamma(\nu_1)} \frac{\Gamma(\nu_2 + \frac{3}{2})}{\Gamma(\nu_2)} \frac{\Gamma(\nu_{12})^2}{\Gamma(\nu_{12} + \frac{3}{2})^2} \frac{a_1^{2\nu_1} a_2^{2\nu_2}}{a_{12}^{4\nu_{12}}} \inf_{t \geq 0} \frac{(a_{12}^2 + t^2)^{2\nu_{12}+3}}{(a_1^2 + t^2)^{\nu_1+(3/2)} (a_2^2 + t^2)^{\nu_2+(3/2)}}. \quad (\text{S.2})$$

Condition (S.2) implies that if $\nu_{12} < \frac{1}{2}(\nu_1 + \nu_2)$, ρ_{12} has to be zero, i.e., the two Cartesian components of \mathbf{Z} are independent. Thus, our primary interest lies in the case where $\nu_{12} \geq \frac{1}{2}(\nu_1 + \nu_2)$.

The derived cross-covariance function of $\mathbf{Y}_{\text{tan},\mathbf{Z}}(\mathbf{s})$ is

$$\begin{aligned} & \mathbf{C}_{\text{tan},\mathbf{Z}}(\mathbf{s}, \mathbf{t}) \\ &= - \begin{pmatrix} \sigma_1 \mathbf{P}_{\mathbf{s}} & \sigma_2 \mathbf{Q}_{\mathbf{s}} \end{pmatrix} \begin{pmatrix} \mathbf{K}(\mathbf{h}; \nu_1, a_1) & \rho_{12} \mathbf{K}(\mathbf{h}; \nu_{12}, a_{12}) \\ \rho_{12} \mathbf{K}(\mathbf{h}; \nu_{12}, a_{12}) & \mathbf{K}(\mathbf{h}; \nu_2, a_2) \end{pmatrix} \begin{pmatrix} \sigma_1 \mathbf{P}_{\mathbf{t}}^{\text{T}} \\ \sigma_2 \mathbf{Q}_{\mathbf{t}}^{\text{T}} \end{pmatrix}, \end{aligned}$$

where $\nu_1, \nu_2 > 1$.

S.1.5 Proof of Proposition 1

The tangential vector field $\mathbf{Y}_{\text{tan},\mathbf{z}}(\mathbf{s})$ with cross-covariance function $\mathbf{C}_{\text{tan},\mathbf{z}}(\mathbf{s}, \mathbf{t})$ can be converted to the canonical coordinates $(\hat{\mathbf{u}}, \hat{\mathbf{v}})$ by the following transformation

$$\mathbf{V}(\mathbf{s}) \equiv (u(\mathbf{s}), v(\mathbf{s}))^T = \mathbf{T}_{\mathbf{s}} \mathbf{Y}_{\text{tan},\mathbf{z}}(\mathbf{s}),$$

where

$$\mathbf{T}_{\mathbf{s}} = \begin{pmatrix} -\sin \phi_{\mathbf{s}} & \cos \phi_{\mathbf{s}} & 0 \\ -\cos \theta_{\mathbf{s}} \cos \phi_{\mathbf{s}} & -\cos \theta_{\mathbf{s}} \sin \phi_{\mathbf{s}} & \sin \theta_{\mathbf{s}} \end{pmatrix}.$$

Then the cross-covariance function of \mathbf{V} is

$$\mathbf{C}_{\mathbf{V}}(\mathbf{s}, \mathbf{t}) = \mathbf{T}_{\mathbf{s}} \mathbf{C}_{\text{tan},\mathbf{z}}(\mathbf{s}, \mathbf{t}) \mathbf{T}_{\mathbf{t}}^T. \quad (\text{S.3})$$

Note that

$$\mathbf{T}_{\mathbf{s}} \mathbf{P}_{\mathbf{s}} = \begin{pmatrix} -\sin \phi_{\mathbf{s}} & \cos \phi_{\mathbf{s}} & 0 \\ -\cos \theta_{\mathbf{s}} \cos \phi_{\mathbf{s}} & -\cos \theta_{\mathbf{s}} \sin \phi_{\mathbf{s}} & \sin \theta_{\mathbf{s}} \end{pmatrix},$$

and

$$\mathbf{T}_{\mathbf{s}} \mathbf{Q}_{\mathbf{s}} = \begin{pmatrix} \cos \theta_{\mathbf{s}} \cos \phi_{\mathbf{s}} & \cos \theta_{\mathbf{s}} \sin \phi_{\mathbf{s}} & -\sin \theta_{\mathbf{s}} \\ -\sin \phi_{\mathbf{s}} & \cos \phi_{\mathbf{s}} & 0 \end{pmatrix}.$$

Plugging them into (S.3) with $\mathbf{s} = \mathbf{t}$, we have

$$\mathbf{C}_{\mathbf{V}}(\mathbf{s}, \mathbf{s}) = -[\sigma_1^2 F_{\text{Mat}}(0; \nu_1, a) + \sigma_2^2 F_{\text{Mat}}(0; \nu_2, a)] \mathbf{I}_2,$$

which is constant with respect to \mathbf{s} .

We can intuitively understand that the co-located cross-correlation between u and v has to be zero through computing it based on an alternative representation of \mathbf{V} , i.e., (8) and (9). Since Z_1 is isotropic, its change rates along latitudinal and longitudinal directions (at

the same location) are uncorrelated (Adler and Taylor, 2009, page 116, (5.7.3)), i.e.,

$$\text{Cov} \left(\frac{\partial Z_1(\theta, \phi)}{\partial \phi}, \frac{\partial Z_1(\theta, \phi)}{\partial \theta} \right) = 0.$$

Similarly, the same holds for Z_2 . Moreover, the cross-covariance between Z_1 and Z_2 is also isotropic. Due to the curvature of the sphere, the covariance between the change rates of Z_1 and Z_2 along latitudinal direction (at the same location) equals the counterpart along longitudinal direction multiplied by $\sin^2 \theta$, i.e.,

$$\text{Cov} \left(\frac{\partial Z_1(\theta, \phi)}{\partial \phi}, \frac{\partial Z_2(\theta, \phi)}{\partial \phi} \right) = \text{Cov} \left(\frac{\partial Z_2(\theta, \phi)}{\partial \theta}, \frac{\partial Z_1(\theta, \phi)}{\partial \theta} \right) \sin^2 \theta.$$

Using the above facts and after certain algebraic manipulations, we arrive at the same conclusion.

S.1.6 Proof of Proposition 2

According to the construction of the TMM, for any (θ, ϕ) except at the two poles,

$$\begin{aligned} u(\theta, \phi) \hat{\mathbf{u}} + v(\theta, \phi) \hat{\mathbf{v}} &= \nabla_{\mathbf{s}}^* \tilde{Z}_1(\mathbf{s}) + L_{\mathbf{s}}^* \tilde{Z}_2(\mathbf{s}) \quad \mathbb{P}\text{-a.e.} \\ &= \left(\frac{1}{\sin \theta} \frac{\partial \tilde{Z}_1}{\partial \phi} + \frac{\partial \tilde{Z}_2}{\partial \theta} \right) \hat{\boldsymbol{\phi}} - \left(\frac{1}{\sin \theta} \frac{\partial \tilde{Z}_2}{\partial \phi} - \frac{\partial \tilde{Z}_1}{\partial \theta} \right) \hat{\boldsymbol{\theta}} \\ &= \left(\frac{1}{\sin \theta} \frac{\partial \tilde{Z}_1}{\partial \phi} + \frac{\partial \tilde{Z}_2}{\partial \theta} \right) \hat{\mathbf{u}} + \left(\frac{1}{\sin \theta} \frac{\partial \tilde{Z}_2}{\partial \phi} - \frac{\partial \tilde{Z}_1}{\partial \theta} \right) \hat{\mathbf{v}}. \end{aligned}$$

Thus,

$$u(\theta, \phi) = \frac{1}{\sin \theta} \frac{\partial \tilde{Z}_1}{\partial \phi} + \frac{\partial \tilde{Z}_2}{\partial \theta} \quad \mathbb{P}\text{-a.e.},$$

and

$$v(\theta, \phi) = \frac{1}{\sin \theta} \frac{\partial \tilde{Z}_2}{\partial \phi} - \frac{\partial \tilde{Z}_1}{\partial \theta} \quad \mathbb{P}\text{-a.e.}$$

By condition **A1** and the chain rule for partial derivatives in quadratic mean (Potthoff, 2010, Theorem 2.13), (8) and (9) follow.

We have assumed that the cross-covariance function of $(Z_1(\mathbf{s}), Z_2(\mathbf{s}))^\top$ (restricted to \mathbb{S}^2) only depends on the chordal distance, which is a function of θ_1, θ_2 and $\phi_1 - \phi_2$. Besides, the coefficients of the partial derivatives are independent of longitude. Using the same argument as Jun and Stein (2008), we conclude that u and v are axially symmetric both marginally and jointly.

S.1.7 Implementation of the DFT for the Cross-covariance Matrix

Suppose that the observations are on a regular latitude-longitude grid on the unit sphere $\{(\theta_i, \phi_j) : i = 1, \dots, n_{\text{lat}}, j = 1, \dots, n_{\text{lon}}\}$, where θ and ϕ represent the co-latitude and longitude, respectively. Besides, $\theta_i - \theta_{i-1} = (\theta_{n_{\text{lat}}} - \theta_1)/(n_{\text{lat}} - 1)$ and $\phi_j - \phi_{j-1} = 2\pi/n_{\text{lon}}$, $\phi_1 = 0$, $\phi_{n_{\text{lon}}} = 2\pi - 2\pi/n_{\text{lon}}$. Let $(V_1(\theta_i, \phi_j), V_2(\theta_i, \phi_j))^\top, i = 1, \dots, n_{\text{lat}}, j = 1, \dots, n_{\text{lon}}$ denote the observations. We rearrange them into $(\mathbf{V}(\theta_1)^\top, \dots, \mathbf{V}(\theta_{n_{\text{lat}}})^\top)^\top$, where

$$\mathbf{V}(\theta_i) = (\mathbf{V}_1(\theta_i)^\top, \mathbf{V}_2(\theta_i)^\top)^\top = (V_1(\theta_i, \phi_1), \dots, V_1(\theta_i, \phi_{n_{\text{lon}}}), V_2(\theta_i, \phi_1), \dots, V_2(\theta_i, \phi_{n_{\text{lon}}}))^\top.$$

Since the cross-covariance function of the random field \mathbf{V} is axially symmetric,

$$\begin{aligned} \text{Cov}(\mathbf{V}(\theta_i), \mathbf{V}(\theta_j)) &= \mathbb{E}(\mathbf{V}(\theta_i)\mathbf{V}(\theta_j)^\top) \\ &= \begin{pmatrix} \mathbb{E}(\mathbf{V}_1(\theta_i)\mathbf{V}_1(\theta_j)^\top) & \mathbb{E}(\mathbf{V}_1(\theta_i)\mathbf{V}_2(\theta_j)^\top) \\ \mathbb{E}(\mathbf{V}_2(\theta_i)\mathbf{V}_1(\theta_j)^\top) & \mathbb{E}(\mathbf{V}_2(\theta_i)\mathbf{V}_2(\theta_j)^\top) \end{pmatrix} \\ &= \begin{pmatrix} [C_{11}(\theta_i, \theta_j, \phi_k - \phi_l)]_{1 \leq k, l \leq n_{\text{lon}}} & [C_{12}(\theta_i, \theta_j, \phi_k - \phi_l)]_{1 \leq k, l \leq n_{\text{lon}}} \\ [C_{21}(\theta_i, \theta_j, \phi_k - \phi_l)]_{1 \leq k, l \leq n_{\text{lon}}} & [C_{22}(\theta_i, \theta_j, \phi_k - \phi_l)]_{1 \leq k, l \leq n_{\text{lon}}} \end{pmatrix} \\ &= \begin{pmatrix} \mathbf{C}_{11}(\theta_i, \theta_j) & \mathbf{C}_{12}(\theta_i, \theta_j) \\ \mathbf{C}_{21}(\theta_i, \theta_j) & \mathbf{C}_{22}(\theta_i, \theta_j) \end{pmatrix}, \end{aligned}$$

where $\mathbf{C}_{..}(\theta_i, \theta_j)$ are circulant matrices of the form (omitting the subscripts $\cdot\cdot$)

$$\begin{pmatrix} c_0 & c_{n_{\text{lon}}-1} & \cdots & c_1 \\ c_1 & c_0 & \cdots & c_2 \\ \vdots & \vdots & & \vdots \\ c_{n_{\text{lon}}-1} & c_{n_{\text{lon}}-2} & \cdots & c_0 \end{pmatrix}.$$

They can be diagonalized by a DFT matrix \mathbf{F} , i.e.,

$$\mathbf{F}\mathbf{C}_{..}(\theta_i, \theta_j)\mathbf{F}^{-1} = \mathbf{\Lambda}_{..},$$

where

$$\mathbf{F} = \begin{pmatrix} 1 & 1 & \cdots & 1 \\ 1 & \omega^1 & \cdots & \omega^{n_{\text{lon}}-1} \\ 1 & \omega^2 & \cdots & \omega^{2(n_{\text{lon}}-1)} \\ \vdots & \vdots & & \vdots \\ 1 & \omega^{n_{\text{lon}}-1} & \cdots & \omega^{(n_{\text{lon}}-1)(n_{\text{lon}}-1)} \end{pmatrix},$$

and $\omega = \exp(-2\pi i/n_{\text{lon}})$. Besides, the diagonal matrix $\mathbf{\Lambda}_{..}$ can be obtained simply by

$$\mathbf{\Lambda}_{..} = \text{diag}(\mathbf{F}\mathbf{c}_{..}),$$

where $\mathbf{c}_{..} = (c_0, \cdots, c_{n_{\text{lon}}-1})^T$. Thus, we transform the observations as follows

$$\begin{aligned} \mathbf{V}(\theta_i) &= (\mathbf{V}_1(\theta_i)^T, \mathbf{V}_2(\theta_i)^T)^T \longrightarrow \\ (\tilde{\mathbf{V}}_1(\theta_i)^T, \tilde{\mathbf{V}}_2(\theta_i)^T)^T &= \left((\mathbf{F}\mathbf{V}_1(\theta_i))^T / \sqrt{n_{\text{lon}}}, (\mathbf{F}\mathbf{V}_2(\theta_i))^T / \sqrt{n_{\text{lon}}} \right)^T, \end{aligned}$$

so that

$$\begin{aligned}
\mathbb{E} \left(\tilde{\mathbf{V}}.(\theta_i) \tilde{\mathbf{V}}.(\theta_j)^{\text{H}} \right) &= \mathbf{F} \mathbb{E} \left(\mathbf{V}.(\theta_i) \mathbf{V}.(\theta_j)^{\text{T}} \right) \mathbf{F}^{\text{H}} / n_{\text{lon}} \\
&= \mathbf{F} \mathbb{E} \left(\mathbf{V}.(\theta_i) \mathbf{V}.(\theta_j)^{\text{T}} \right) \mathbf{F}^{-1} \\
&= \mathbf{\Lambda}.,
\end{aligned}$$

where \cdot^{H} represents the Hermitian transpose. We further rearrange the transformed observations into

$$\begin{aligned}
&\left(\tilde{V}_1(\theta_1, \phi_1), \dots, \tilde{V}_1(\theta_{n_{\text{lat}}}, \phi_1), \dots, \tilde{V}_1(\theta_1, \phi_{n_{\text{lon}}}), \dots, \tilde{V}_1(\theta_{n_{\text{lat}}}, \phi_{n_{\text{lon}}}), \right. \\
&\quad \left. \tilde{V}_2(\theta_1, \phi_1), \dots, \tilde{V}_2(\theta_{n_{\text{lat}}}, \phi_1), \dots, \tilde{V}_2(\theta_1, \phi_{n_{\text{lon}}}), \dots, \tilde{V}_2(\theta_{n_{\text{lat}}}, \phi_{n_{\text{lon}}}) \right)^{\text{T}}.
\end{aligned}$$

Then the corresponding cross-covariance matrix has the form

$$\mathbf{\Sigma} = \begin{pmatrix} \mathbf{\Sigma}_1 & \mathbf{\Sigma}_{12} \\ \mathbf{\Sigma}_{12}^{\text{H}} & \mathbf{\Sigma}_2 \end{pmatrix},$$

where $\mathbf{\Sigma}_1$, $\mathbf{\Sigma}_{12}$ and $\mathbf{\Sigma}_2$ are complex block diagonal matrices. We know that the determinant of $\mathbf{\Sigma}$

$$\det(\mathbf{\Sigma}) = \det(\mathbf{\Sigma}_1 - \mathbf{\Sigma}_{12} \mathbf{\Sigma}_2^{-1} \mathbf{\Sigma}_{12}^{\text{H}}) \det(\mathbf{\Sigma}_2),$$

and the inverse of $\mathbf{\Sigma}$

$$\mathbf{\Sigma}^{-1} = \begin{pmatrix} (\mathbf{\Sigma}_1 - \mathbf{\Sigma}_{12} \mathbf{\Sigma}_2^{-1} \mathbf{\Sigma}_{12}^{\text{H}})^{-1} & -\mathbf{\Sigma}_1^{-1} \mathbf{\Sigma}_{12} (\mathbf{\Sigma}_2 - \mathbf{\Sigma}_{12}^{\text{H}} \mathbf{\Sigma}_1^{-1} \mathbf{\Sigma}_{12})^{-1} \\ (-\mathbf{\Sigma}_1^{-1} \mathbf{\Sigma}_{12} (\mathbf{\Sigma}_2 - \mathbf{\Sigma}_{12}^{\text{H}} \mathbf{\Sigma}_1^{-1} \mathbf{\Sigma}_{12})^{-1})^{\text{H}} & (\mathbf{\Sigma}_2 - \mathbf{\Sigma}_{12}^{\text{H}} \mathbf{\Sigma}_1^{-1} \mathbf{\Sigma}_{12})^{-1} \end{pmatrix}.$$

Using these two facts, all the matrix computations can be reduced to those of block diagonal matrices, which implies that both $\det(\mathbf{\Sigma})$ and $\mathbf{\Sigma}^{-1}$ can be computed efficiently.

S.1.8 Initial Value Specification

To specify initial values for the model parameters, we first fix the spatial scale parameter at a sufficiently large value (e.g., 5 in the simulation study and 10 in the real data analysis). This is motivated by the numerical results in [Kaufman and Shaby \(2013\)](#) that when the spatial scale parameter is fixed at some value larger than its true value, both the estimators and predictors can still perform well. We set initial values for τ_1 and τ_2 at certain relatively small values since they have little impact on the parameter estimates. For the remaining parameters, we randomly select 100 parameter vectors in the parameter space using Latin hypercube sampling (LHS), and choose the one with the smallest negative log-likelihood.

S.1.9 Computational Details of the NBG

This subsection describes the computational details of the NBG. Note that currently there is no publicly available package for it. The model consists of two components: one is linear combinations of partial derivatives of a scalar random field, and the other is a PARS-BM with $\rho_{12} = 0$. This special formulation enables us to obtain rough estimates of the parameters, which can be served as initial values for another algorithm for further refinement. Specifically, one can start with estimating the parameters in the second component by ignoring those in the first one (i.e., assuming $a_i = b_i = 0$), and then conditionally on the parameter estimates just obtained, estimate the remaining parameters in the first component. The estimates of all the parameters can be used as initial values in the joint minimization of the negative log-likelihood function. In all cases except the estimation procedure for computing bootstrap standard errors, we use the *interior-point* algorithm through the Matlab function *fmincon* to jointly minimize the negative log-likelihood function. The maximum number of function evaluations for the algorithm is specified as 20,000 to estimate the parameters in the real data analysis, and 10,000 otherwise. Note that the computation of bootstrap standard errors is time-consuming since the estimation procedure needs to be repeated 200 times. To reduce the computational burden with a mild loss of accuracy, instead we use the *Nelder-Mead*

simplex direct search algorithm through the Matlab function *fminsearch* in the estimation procedure. Although this is an unconstrained optimization method, the results will not be significantly affected as long as initial values are set close to the actual minimizer.

S.2 Additional Numerical Results

S.2.1 Accuracy of Parameter Estimation on Irregular Grids

To investigate the accuracy of parameter estimation on irregular grids, we conduct some smaller-scale experiments (due to the computational burden). Specifically, we generate 100 realizations with

$$\boldsymbol{\theta} = (1, 1, 0.5, 3, 4, 1/2, 0.1, 0.1)$$

on two HEALPix grids (Górski et al., 2005) with 192 and 768 grid points, respectively. The HEALPix grid partitions the unit sphere into equal area pixels, and hence is naturally irregular. The results are summarized by boxplots in Figure S.1.

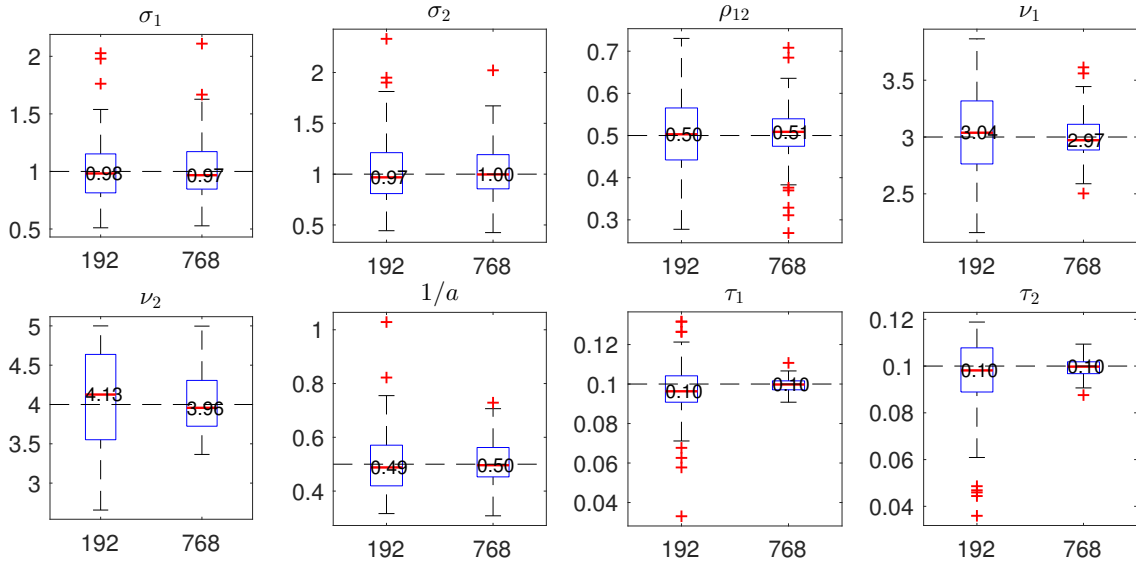


Figure S.1: Results of the Monte Carlo simulation study to investigate the accuracy of parameter estimation for the TMM on irregular grids. The MLEs of $\boldsymbol{\theta} = (\sigma_1, \sigma_2, \rho_{12}, \nu_1, \nu_2, 1/a, \tau_1, \tau_2)$ are summarized by boxplots for two simulated data sets with increasing sample size (shown on x-axis). The dashed horizontal lines are at the true values.

S.2.2 Accuracy of Parameter Estimation For Three Different Noise Levels

It is also interesting to see how the noise level affects the accuracy of parameter estimation. Specifically, we consider three different noise levels, and generate 500 realizations with

$$\boldsymbol{\theta} = (1, 1, 0.5, 3, 4, 1/2, 0.1, 0.1) \quad \text{low noise}$$

$$\boldsymbol{\theta} = (1, 1, 0.5, 3, 4, 1/2, 0.2, 0.2) \quad \text{medium noise}$$

$$\boldsymbol{\theta} = (1, 1, 0.5, 3, 4, 1/2, 0.3, 0.3) \quad \text{high noise}$$

on a regular grid with $(n_{\text{lat}}, n_{\text{lon}}) = (15, 30)$, respectively. Figure S.2 shows that the higher the noise level, the larger the spreads of the estimates, especially for ν_1 and ν_2 .

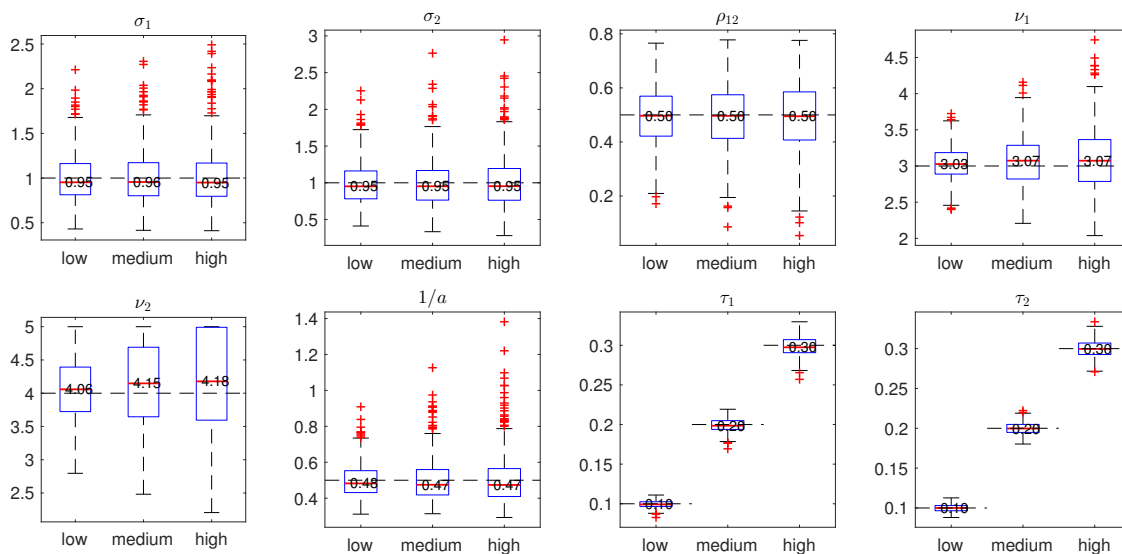


Figure S.2: Results of the Monte Carlo simulation study to investigate the accuracy of parameter estimation for the TMM for three different noise levels (low: $\tau_1 = \tau_2 = 0.1$; medium: $\tau_1 = \tau_2 = 0.2$; high: $\tau_1 = \tau_2 = 0.3$). The MLEs of $\boldsymbol{\theta} = (\sigma_1, \sigma_2, \rho_{12}, \nu_1, \nu_2, 1/a, \tau_1, \tau_2)$ are summarized by boxplots. The dashed horizontal lines are at the true values.

S.2.3 Accuracy of Parameter Estimation For Two Different Sets of Smoothness Parameters

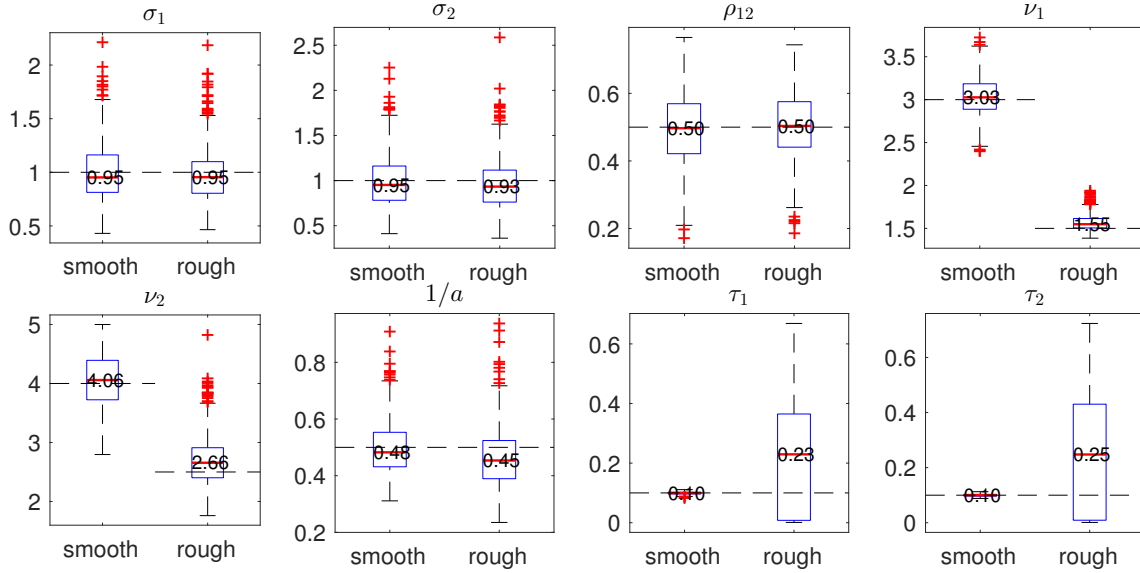


Figure S.3: Results of the Monte Carlo simulation study to investigate the accuracy of parameter estimation for the TMM for two different sets of smoothness parameters (smooth: $\nu_1 = 3, \nu_2 = 4$; rough: $\nu_1 = 1.5, \nu_2 = 2.5$). The MLEs of $\boldsymbol{\theta} = (\sigma_1, \sigma_2, \rho_{12}, \nu_1, \nu_2, 1/a, \tau_1, \tau_2)$ are summarized by boxplots. The dashed horizontal lines are at the true values.

We also test the effect of the smoothness of the vector field on the accuracy of parameter estimation. Specifically, we consider two different sets of smoothness parameters, and generate 500 realizations with

$$\boldsymbol{\theta} = (1, 1, 0.5, 3, 4, 1/2, 0.1, 0.1) \quad \text{smooth}$$

$$\boldsymbol{\theta} = (1, 1, 0.5, 1.5, 2.5, 1/2, 0.1, 0.1) \quad \text{rough}$$

on a regular grid with $(n_{\text{lat}}, n_{\text{lon}}) = (15, 30)$, respectively.

The results are summarized by boxplots in Figure S.3. Due to the application of the differential operators, the latter specification yields a vector field effectively as rough as that generated by the Matérn model with $\nu = 1.5 - 1 = 0.5$. We know that it is difficult to distinguish between a very rough vector field and observational (white) noises. Thus, in the

second case, τ_1 and τ_2 are significantly overestimated, and their estimates have large variability. Moreover, compared with those in the first case, the upward biases of the estimates for ν_1 and ν_2 in the second case are more pronounced.

S.2.4 Accuracy of Parameter Estimation When Covariates Are Included in the Model

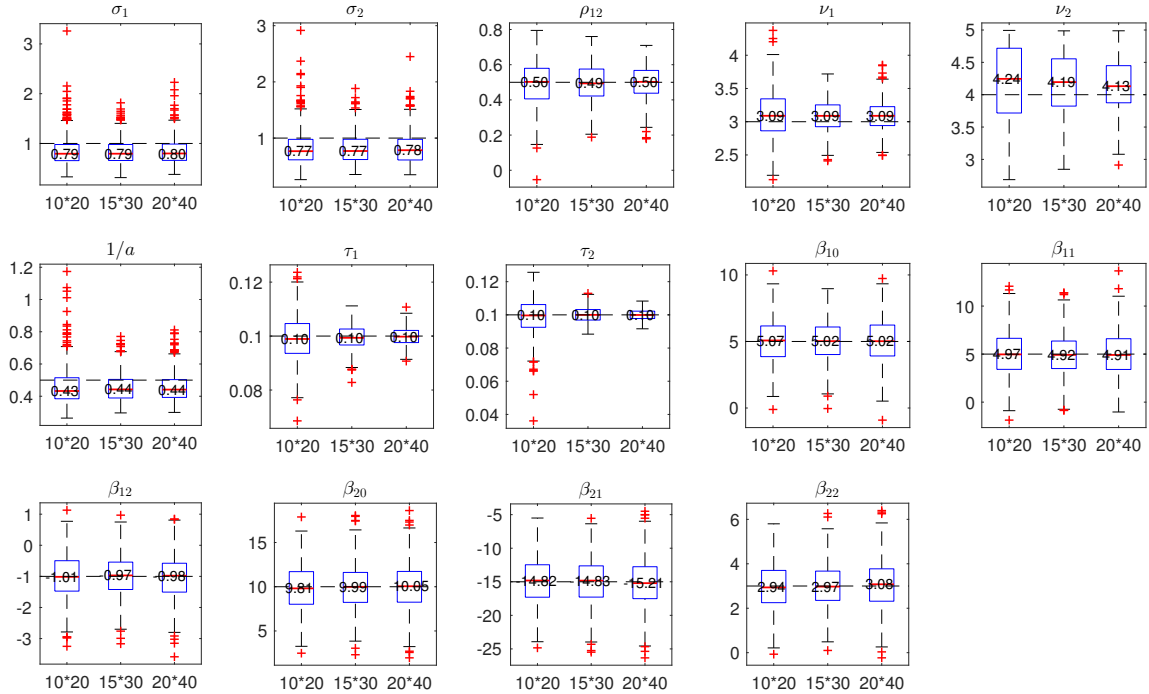


Figure S.4: Results of the Monte Carlo simulation study to investigate the accuracy of parameter estimation for the TMM when covariates are included in the model. The MLEs of $\boldsymbol{\theta} = (\sigma_1, \sigma_2, \rho_{12}, \nu_1, \nu_2, 1/a, \tau_1, \tau_2)$ and $\boldsymbol{\beta}_i, i = 0, 1, 2$ are summarized by boxplots for three simulated data sets with increasing sample size (shown on x-axis). The dashed horizontal lines are at the true values.

We include a polynomial function of the co-latitude as covariates in the model such that for the vector field $\mathbf{Y}(\mathbf{s}), \mathbf{s} \equiv (\theta, \phi)$, its mean is $\boldsymbol{\beta}_0 + \boldsymbol{\beta}_1\theta + \boldsymbol{\beta}_2\theta^2$, where $\boldsymbol{\beta}_i = (\beta_{1i}, \beta_{2i})^\top, i = 0, 1, 2$ are two-dimensional coefficient vectors, and its cross-covariance function still follows

the TMM (augmented with nugget effects). We generate 500 realizations with

$$\boldsymbol{\theta} = (1, 1, 0.5, 3, 4, 1/2, 0.1, 0.1),$$

and

$$\boldsymbol{\beta}_0 = (5, 10)^T, \boldsymbol{\beta}_1 = (5, -15)^T, \boldsymbol{\beta}_2 = (-1, 3)^T$$

on three regular grids with $(n_{\text{lat}}, n_{\text{lon}}) = (10, 20), (15, 30)$ and $(20, 40)$, respectively.

Figure S.4 shows that the estimates for σ_1 , σ_2 and $1/a$ have larger biases than those obtained without any covariate. This suggests difficulty in jointly estimating the parameters in both the mean and the random components. The biases of the estimates for $\boldsymbol{\beta}_i, i = 0, 1, 2$ are small, while the spreads of the estimates do not decrease significantly as the sample size increases.

S.2.5 Bootstrap Standard Errors

The asymptotic standard errors of the MLEs based on the Fisher information matrix are not suitable in our context since their validity depends on the framework of increasing domain asymptotics. Therefore, we seek to estimate the standard errors of the MLEs using a parametric bootstrap procedure. Specifically, the bootstrap samples are independent realizations of the fitted TMM with its parameters estimated by the method of maximum likelihood (Horowitz, 2001). To examine the effectiveness of the parametric bootstrap procedure, using the simulation study in Section 3.2 with $(n_{\text{lat}}, n_{\text{lon}}) = (15, 30)$ as an example, we compare the empirical standard errors, and the bootstrap standard errors computed based on 200 bootstrap samples. The empirical standard errors, which can be regarded as estimates for the true standard errors, are computed as the standard errors of the estimates obtained from the 500 simulation runs. We obtain 10 sets of bootstrap standard errors for the first 10 of the 500 simulation runs. Figure S.5 displays the ratios between the bootstrap and empirical standard errors. For each parameter, the 10 ratios are summarized by a boxplot. Note

that the bootstrap standard errors tend to be slightly larger than the empirical ones, which suggests that the former can be seen as conservative estimates for the true standard errors. Also, the ratios for σ_1 and σ_2 have relatively large spreads.

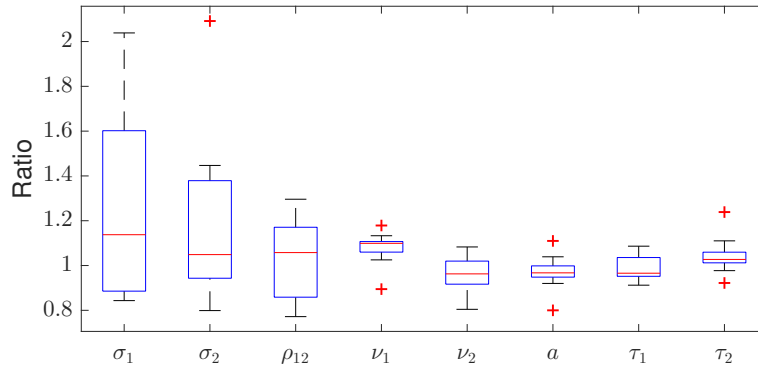


Figure S.5: Ratios between the bootstrap and empirical standard errors. For each parameter, the 10 ratios are summarized by a boxplot.

S.2.6 Predictive Performance Comparison

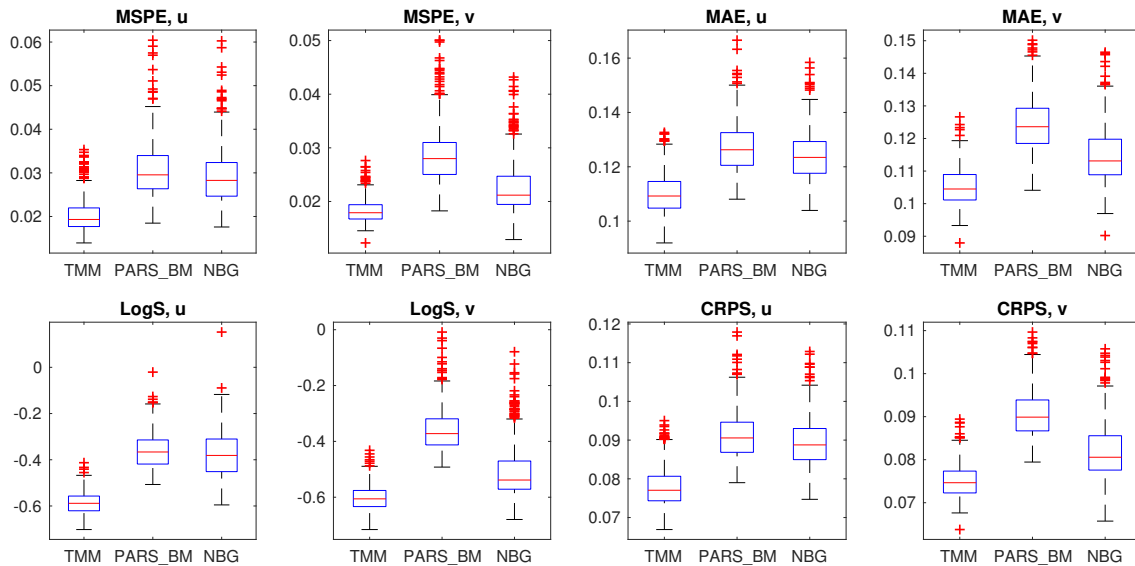


Figure S.6: Boxplots of MSPE, MAE, LogS and CRPS for the TMM, PARS-BM and NBG with 500 replications.

In this subsection, we compare the predictive performance of the TMM, PARS-BM and NBG by a Monte Carlo simulation study. The spatial prediction of vector fields can be

performed by cokriging; for example, Myers (1982); Ver Hoef and Cressie (1993). We fit the PARS-BM using the R package RandomFields (Version 3.0.62) (Schlather et al., 2015), and the computational details of the NBG are given in Supplementary Materials S.1.9. Suppose that the data are simulated from the TMM with $\boldsymbol{\theta} = (1, 1, 0.5, 3, 4, 1/2, 0.1, 0.1)$ on the HEALPix grid with 768 grid points. Half of the locations are randomly selected outside a randomly selected longitudinal region with width 30° for model fitting, and the remaining locations are held out for spatial prediction. Both short-range and long-range predictions are considered through separating the data in this way. We assess the prediction accuracy using four popular scoring rules: the mean squared prediction error (MSPE), the mean absolute error (MAE), the logarithmic score (LogS), and the continuous ranked probability score (CRPS) (Gneiting and Raftery, 2007). The cross-validation procedure is repeated 500 times, and the four scoring rules for the three models are summarized by boxplots in Figure S.6. All the boxplots for the TMM are the lowest, which indicates that it has the best predictive performance in terms of the scoring rules.

S.3 Additional Data Example Results

S.3.1 Temporal Independence and Normality Check

Figures S.7 and S.8 display the sample autocorrelation functions (over time) of the u and v residual wind fields, respectively, at 16 randomly selected locations in the subregion of the Indian Ocean. They support our assumption that the residuals are temporally uncorrelated.

The marginal Q-Q plots and the Chi-Square Q-Q plot, shown in Figure S.9 for January, 2000, suggest that the Gaussianity assumption on the u and v residual wind fields is reasonable.

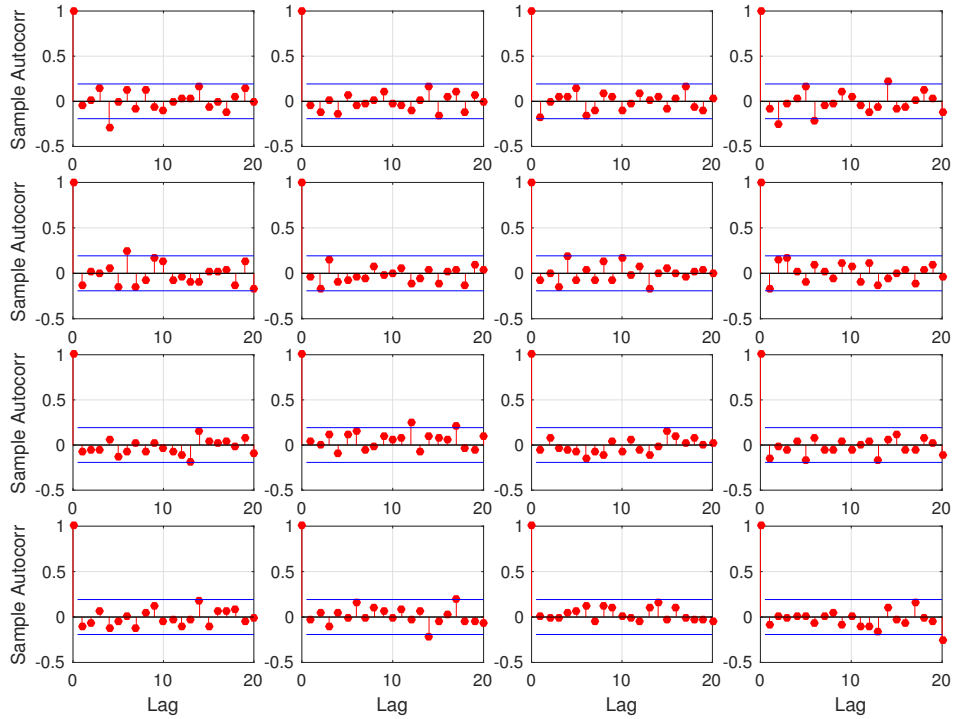


Figure S.7: Sample autocorrelation functions (over time) of the u residual wind fields at 16 randomly selected locations in the subregion of the Indian Ocean.

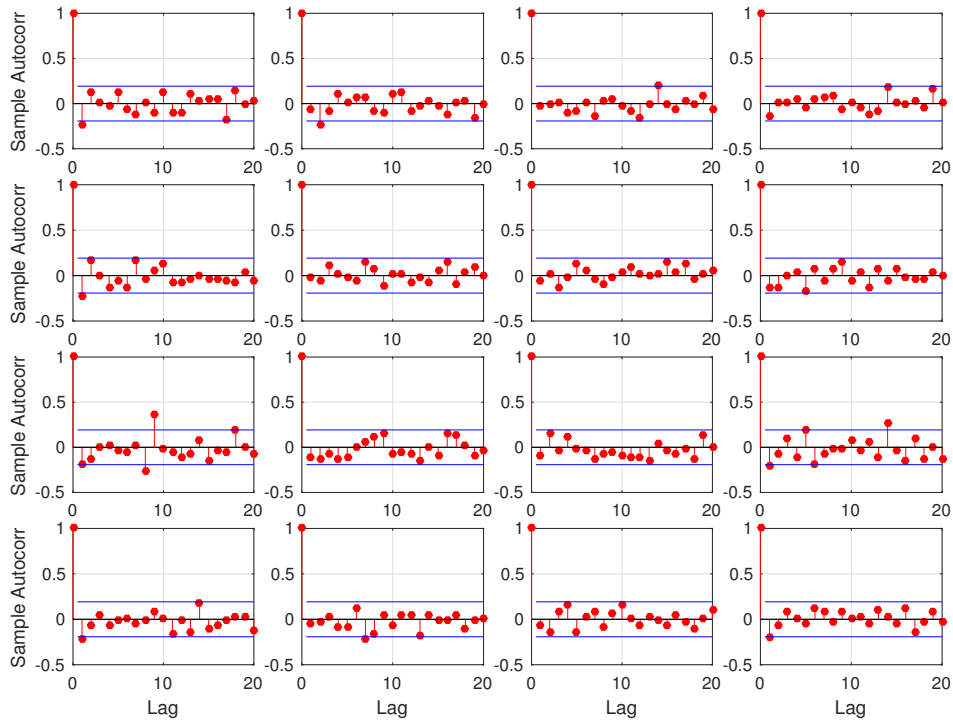


Figure S.8: Sample autocorrelation functions (over time) of the v residual wind fields at 16 randomly selected locations in the subregion of the Indian Ocean.

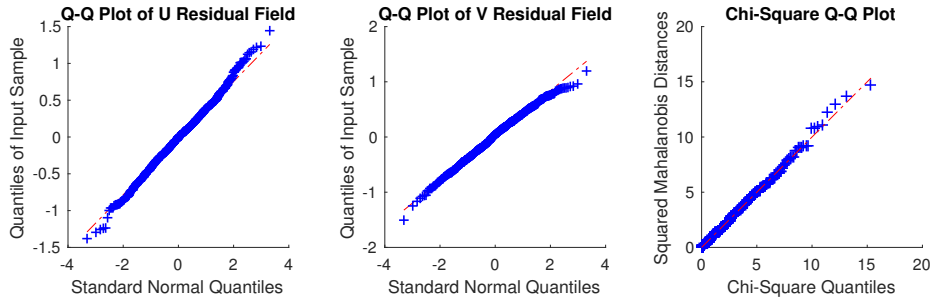


Figure S.9: Marginal Q-Q plots (**the first two**) and Chi-Square Q-Q plot (**the third**) of the u and v residual wind fields for January, 2000 in the subregion of the Indian Ocean.

S.3.2 Comparison between the Empirical and Fitted Variances of the u and v Residual Wind Fields

Figure S.10 displays the empirical and fitted variances of the u and v residual wind fields under the TMM, PARS-BM and NBG, which are plotted with respect to latitude and longitude. The fitted variances of the u and v residual wind fields under the TMM and PARS-BM are constant due to Proposition 1 and the isotropy of the model, respectively. Although they are supposed to be non-stationary with respect to latitude, the fitted variances under the NBG seem nearly identical to those under the other two models.

The empirical variances of the u residual wind fields show a clear pattern depending on latitude, which leads to a noticeable discrepancy between the empirical and fitted variances. This is not surprising because all the three models are too simple to capture this pattern given the number of parameters they have. As discussed in Section 5, the discrepancy for the TMM can be potentially alleviated by assuming a more complicated underlying potential field, such as an anisotropic one. Moreover, one may use other scientific models instead of VEOFs to represent the large-scale component so that the pattern is removed from the residuals.

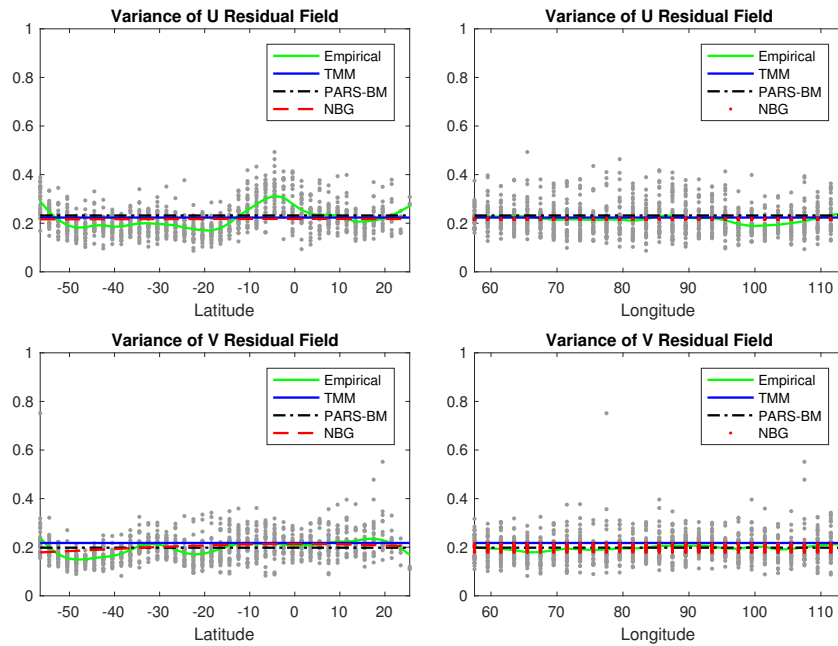


Figure S.10: Empirical and fitted variances of the u and v residual wind fields: circles for the empirical variances, green solid lines for the loess curves fitted to the empirical variances, blue solid lines for the TMM, dash-dot lines for the PARS-BM, and dashed lines (**left**) or dots (**right**) for the NBG. They are plotted with respect to latitude (**left**) and longitude (**right**).

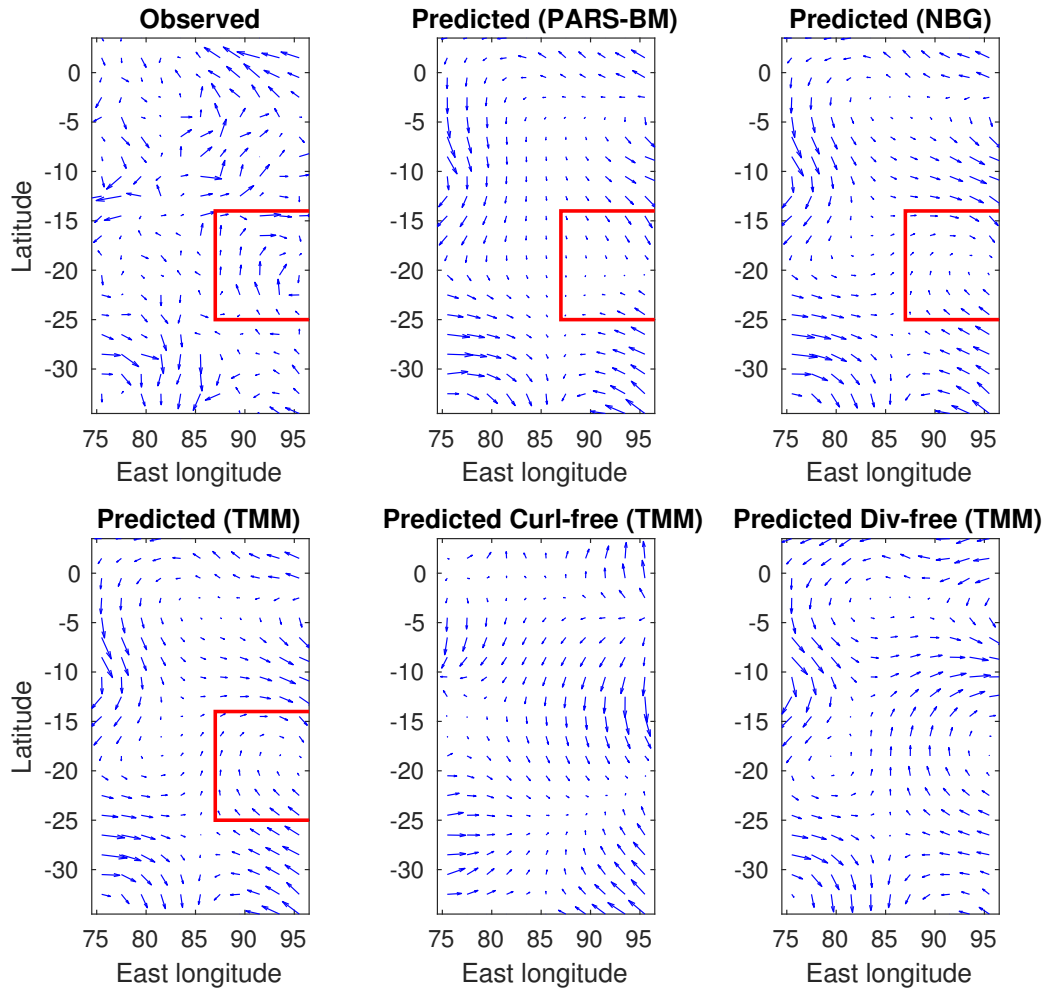


Figure S.11: Observed and predicted wind fields by the TMM, PARS-BM and NBG for January, 2000 in the rectangular region in the center of the subregion of the Indian Ocean. The wind field predicted by the TMM is decomposed into its curl-free and divergence-free components. Note that the arrows have been automatically scaled to fit within the grid.

S.3.3 Predicted and Observed Wind Field Comparison

In this subsection, we compare the wind fields predicted by the TMM, PARS-BM and NBG with the observations. The predicted wind fields are obtained by cokriging, which predicts at the locations in the rectangular region based on the observations outside the region, using the parameter estimates from the entire data set (see Table 1). Figure S.11 displays the observed and predicted wind fields for January, 2000 in the rectangular region, together with a decomposition of the wind field predicted by the TMM into its curl-free and divergence-free components. It is worth pointing out that the wind fields predicted by the TMM and NBG are in good agreement with the observations in terms of certain key physical characteristics, which are missed in the wind field predicted by the PARS-BM. Specifically, in the region marked in red, the observed and predicted winds by the TMM and NBG all rotate clockwise and show similar patterns.

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